## Conical emission from DC-biased filament at 10 THz

N.A. Panov<sup>1,2</sup>, G.E. Rizaev<sup>2</sup>, D.E. Shipilo<sup>1,2</sup>, D.V. Pushkarev<sup>1,2</sup>, I.A. Nikolaeva<sup>1,2</sup>, D.V. Mokrousova<sup>1,2</sup>, L.V. Seleznev<sup>1,2</sup>, A.A. Ionin<sup>2</sup>, O.G. Kosareva<sup>1,2</sup>

1- Faculty of Physics, Lomonosov Moscow State University, 119991, Moscow, Russia 2- P.N. Lebedev Physical Institute of the RAS, 119991, Moscow, Russia napanov@ilc.edu.ru

Femtosecond filament in the external electrostatic field (DC-biased filament) is a prominent source of terahertz (THz) radiation [1]. As the external field grows above  $\sim 3$  kV/cm [2], the directional diagram of the THz emission from a DC-biased filament in air is unimodal with the maximum on the laser beam axis. This was confirmed by numerous experiments and numerous registration techniques, either narrowband detection [2], wideband detection [3, 4], or spectrally resolved [5, 6]. The excellent directionality of THz emission from DC-biased filament makes this THz source a promising tool for the measurements of low  $(10^{13}-10^{15} \text{ cm}^{-3})$  densities of free electrons [7]. The measurements of THz directional diagrams [2–6] were done in the low-frequency range below 2–3 THz. However, 3D + time simulations of THz generation in DC-biased filament performed in our recent work [6] predicted the appearance of THz conical emission in high-frequency THz range (for our 90-fs pulse at  $\sim 10$  THz). In this work, we confirm this prediction experimentally.

In our experiment, we focused the 740-nm, 90-fs, 1.8-mJ pulse into the air gap between the electrodes biased by 15-kV/cm static electric field. The plasma filament between the electrodes was a source of THz radiation detected by a superconducting MoRe bolometer Scontel RS-CCR-1-12T-1+0.3-3T-0.1 sensitive in the spectral range of 0.3-10 THz. The bolometer with the bandpass filters (centered at the frequencies  $\nu = 0.5$ , 1, 3 and 10 THz) was fixed on the 40-cm-long horizontal board and rotated at the horizontal angle  $\alpha$  around the vertical axis. The spherical mirror was fixed on the vertical 40-cm post. To vary the vertical angle  $\beta$ , we moved the focusing mirror along the post. So, the variation of the angles  $\alpha$  and  $\beta$  allowed us to reconstruct the 2D distributions of the THz fluence  $F(\alpha, \beta)$  at the frequency  $\nu$  determined by the bandpass filter. We traced experimentally the transit from the on-axis unimodal angular distribution at 0.5-1 THz to the conical one at 10 THz, see Fig. 1.

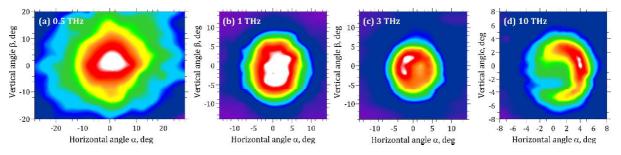


Fig. 1. 2D distributions of the fluence  $F(\alpha, \beta)$  measured at frequencies (a)  $\nu = 0.5$  THz, (b) 1 THz, (c) 3 THz, (d) 10 THz.

- [1] T. Löffler, F. Jacob, and H. Roskos, Appl. Phys. Lett. 77, 453 (2000).
- [2] A. Houard, Y. Liu, B. Prade, V. T. Tikhonchuk, and A. Mysyrowicz, Phys. Rev. Lett. 100, 255006 (2008).
- [3] R. Akhmedzhanov, I. Ilyakov, V. Mironov, E. Suvorov, D. Fadeev, and B. Shishkin, Radiophys. Quantum Electron. 52, 482 (2009).
- [4] D. Lubenko, V. Prokopev, S. Alekseev, M. Ivanov, and V. Losev, Atmospheric Ocean. Opt. 32, 430 (2019).
- [5] T. Fukuda, T. Otsuka, Y. Sentoku, H. Nagatomo, H. Sakagami, R. Kodama, and N. Yugami, Jpn. J. Appl. Phys. 59, 020902 (2020).
- [6] I. Nikolaeva, D. Shipilo, D. Pushkarev, G. Rizaev, D. Mokrousova, A. Koribut, Y. Grudtsyn, N. Panov, L. Seleznev, W. Liu, A. Ionin, and O. Kosareva, Opt. Lett. 46, 5497 (2021).
- [7] D. Jang, H. S. Uhm, D. Jang, M. S. Hur, and H. Suk, Plasma Sources Sci. Technol. 25, 065008 (2016).